

Simulation of Phase Transition and its Visualization



2024/12/15



Basic Concept

Method

CONTANTS

Result discussion

Conclusion



Part I Basic Concept of Phase transition





Phase Transition

• **Classical Phase Transition**: caused by thermal parameters e.g. water-ice-vapor, superconductor-resistor, <u>Ising Model</u>

• Quantum Phase Transition: adulteration, magnetization e.g. Bose-Hubbard Model

Basic Concept

Method



2D Ising Model

- Paramagnetic-ferromagnetic, T_c : Curie Temperature
- Lenz Ising : 1D
- Mean field theory: $T_c = 2$
- Ensemble Theory, Onsager, Second-order transition, $\frac{d^2E}{dT^2}$ discontinued
- Landau: Order parameter symmetry breaking theory





Part II Method



- 1. Give a random spin configuration $\{s_i\} \equiv s_1, s_2 \cdots, s_N$
- 2. Select a point randomly, calculate the energy change if flip it.

$$\Delta E = (-J(-s_m)\sum_n s_n) - (-Js_m\sum_n s_n) = 2Js_m\sum_n s_n$$

- 3. If $\Delta E \leq 0$, flip; if $\Delta E > 0$, flip it by a ratio $A = e^{-\Delta E/kT} \in [0, 1]$
- 4. Repeat 2-3, (Markov chain length)
- 5. Change the Temperature and repeat 2-4.



Important sample



Markov process

- 1. Ergodicity 🗹
- 2. Detailed Balance:

$$p_{\mu} P(\mu \to \nu) = p_{\nu} P(\nu \to \mu)$$
$$p_{\mu} = e^{-\beta E_{\mu}}, \qquad p_{\nu} = e^{-\beta E_{\nu}}$$
$$P(\mu \to \nu) \equiv g(\mu \to \nu) \cdot A(\mu \to \nu)$$
$$g_{\mu \to \nu} = \frac{1}{N}$$

$$\frac{P(\mu \to \nu)}{P(\nu \to \mu)} = \frac{g(\mu \to \nu)A(\mu \to \nu)}{g(\nu \to \mu)A(\nu \to \mu)} = \frac{A(\mu \to \nu)}{A(\nu \to \mu)} = e^{-\beta (E_{\nu} - E_{\mu})}$$

 $A(\mu \rightarrow \nu) = \begin{cases} e^{-\beta (E_{\nu} - E_{\mu})}, & \text{if } E_{\nu} - E_{\mu} > 0, \\ 1, & \text{others.} \end{cases}$





Metropolis's Shortcoming

- 1. Single-Filp-Dynamics: convergence speed is too slow.
- 2. Easy to fall into metastable state.



Wolff Algorithm





Wolff's process

- 1. Give a random spin configuration $\{s_i\} \equiv s_1, s_2 \cdots, s_N$
- 2. Select a seed point randomly,
- 3. Grow the cluster: consider its neighbors, add them to the cluster with $P_{add} = 1 e^{2J/kT}$ (use random numbers), until no new members.
- 4. Flip the cluster.
- 5. Repeat step 2-4. (called Wolff length)
- 6. Change *T*, repeat 2-5.

Basic Concept Method

Conclusion



c: number of added ,

m: number of not added at μ ,

n: number of not added at ν .

$$g(\mu \to \nu) = (P_{add})^{c}(1 - P_{add})^{m},$$

$$g(\nu \to \mu) = (P_{add})^{c}(1 - P_{add})^{n}$$

$$e^{-\beta (E_{\nu} - E_{\mu})} = \frac{P(\mu \to \nu)}{P(\nu \to \mu)} = \frac{g(\mu \to \nu)A(\mu \to \nu)}{g(\nu \to \mu)A(\nu \to \mu)} = (1 - P_{add})^{m-n} \frac{A(\mu \to \nu)}{A(\nu \to \mu)}$$

$$\Delta E = E_{\nu} - E_{\mu} = 2J(m-n)$$

$$\Rightarrow \frac{A(\mu \to \nu)}{A(\nu \to \mu)} = e^{\beta 2J(m-n)}(1 - P_{add})^{m-n}$$
Wolff: $P_{add} = 1 - e^{-2\beta J}$

$$\Rightarrow A_{w}(\nu \to \mu) = A_{w}(\mu \to \nu) = 1$$

- Accepted ratio: Wolff > Metropolis $A_m(\mu \rightarrow \nu) = \begin{cases} e^{-\beta (E_\nu E_\mu)}, & \text{if } E_\nu E_\mu > 0, \\ 1, & \text{others.} \end{cases}$
- A sample process will definitely bring a flip: A higher efficiency.

Method

Result & Discussion

Conclusion

 $\Delta E = 2 J s_m \sum s_n$







Conclusion



Periodical Boundary Condition





Part III Result and Discussion

Conclusion



1. Order parameter \bar{s} -- the best result



Basic Concept	Method	Result & Discussion	Conclusion	

Reduce/enlarge the system's size: Metropolis is fast but worsen.



Wolff, L = 100, $l_w = 1000$, t = 690 s



Metorpolis, L = 400, $l_m = 500$ k, t = 159 s

Metorpolis, L = 40, $l_m = 500$ k, t = 44 s



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Wolff, L = 40, $l_w = 1000$, t = 109 s

Metropolis

3

3.5

- Analytical



- Enlarge the size,
 - Wolff have a higher precision (sacrifice time)
 - Metropolis could not find all of the states, faster but larger errors.

If approaching to the Thermodynamics limit:
 Wolff ~ Idea Ising Model.



2. Energy & Specific heat capacity 2th phase transition



Wolff $L = 100, l_w = 2000, t = 190 \text{ s}$

Metropolis $L = 100, l_m = 500$ k, t = 108s





T = 1, L = 40,100,400, t = 44 s



T = 3, L = 40, 100, 400, t = 46 s



- Convergence speed at low/high T
 - Metropolis is slower, need more sampling steps.
 In larger size it has a probability of unconverging.

Basic Concept	Method	Result & Discussion	Conclusion		
				ANAW 1923 SS	
4 Near T					

		表 1 $T_c = 2.26$ 下两种算法的误差对比							
		1	2	3	4	5	平均值	标准差	相对误差
Analytical	Metropolis	0.7843	0.7469	0.7412	0.7776	0.7719	0.7644	0.0192	0.31%
$\bar{s} = 0.6134$	Wolff	0.6120	0.6207	0.6161	0.6159	0.6119	0.6153	0.0036	24.6%

Wolff's result is more stable and more precise

• Metropolis's single-flip strategy is hard to simulate all states of the system.



Part IV Conclusion



Metropolis, low efficiency, needs more samples (longer markov chains) ,

imprecise, can not improved by only increase the size.

just use as an approximation.

In certain size it is acceptable.

• Wolff is faster in high temperature. Near critical point, the result is stable because of finding out more possible state.





Visualization of Ising model's phase transition process





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